

Spacetime-emergent ring toward tabletop quantum gravity experiments

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Based on arXiv: 2211.13863

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@ Osaka University

Holographic quantum gravity experiment?

True QG theory?

We need QG experiments

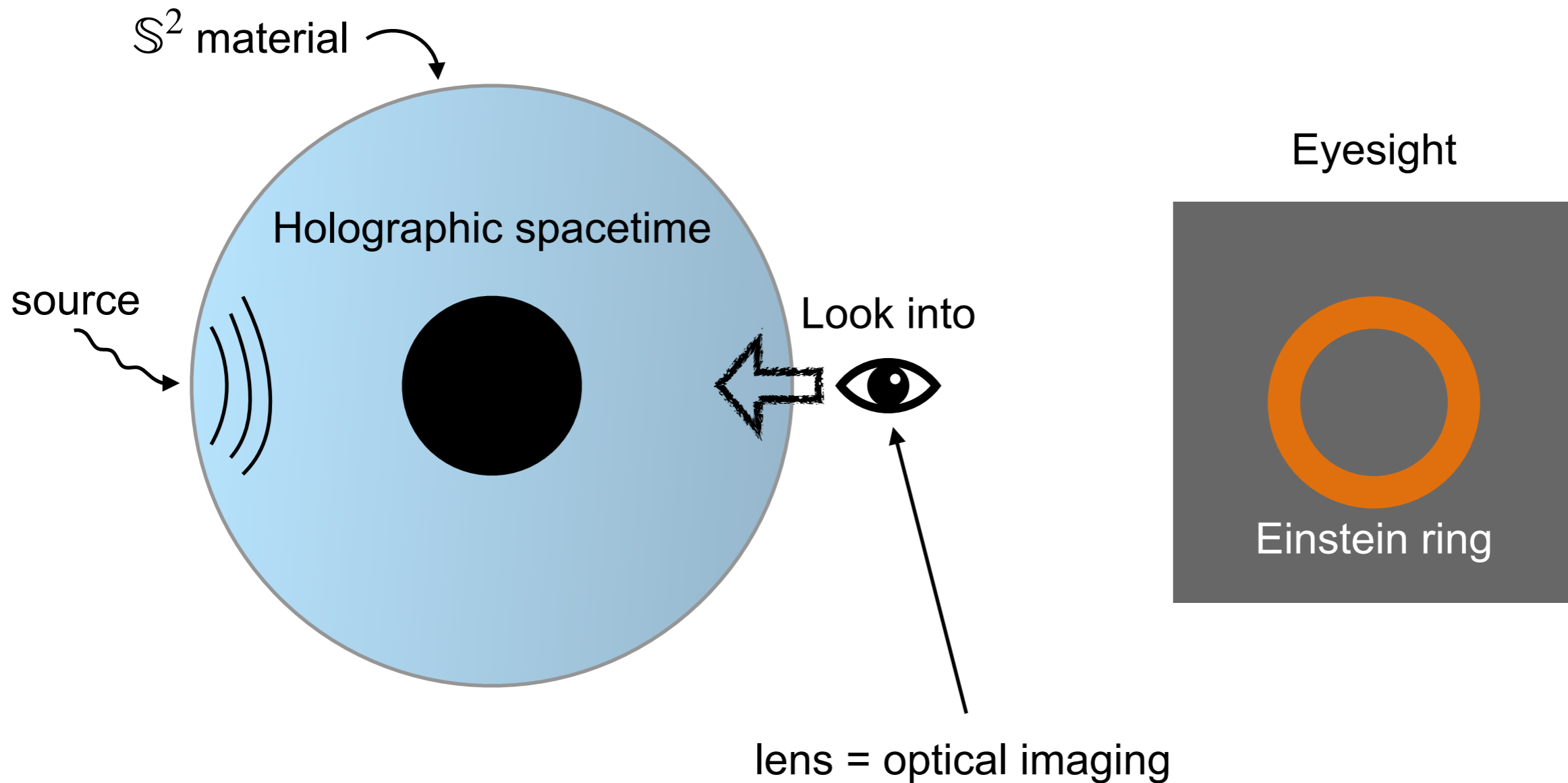
One direction: AdS/CFT

Find materials following the AdS/CFT

How to find?

Seeing the holographic spacetime

Hashimoto, Kinoshita, Murata (2018)

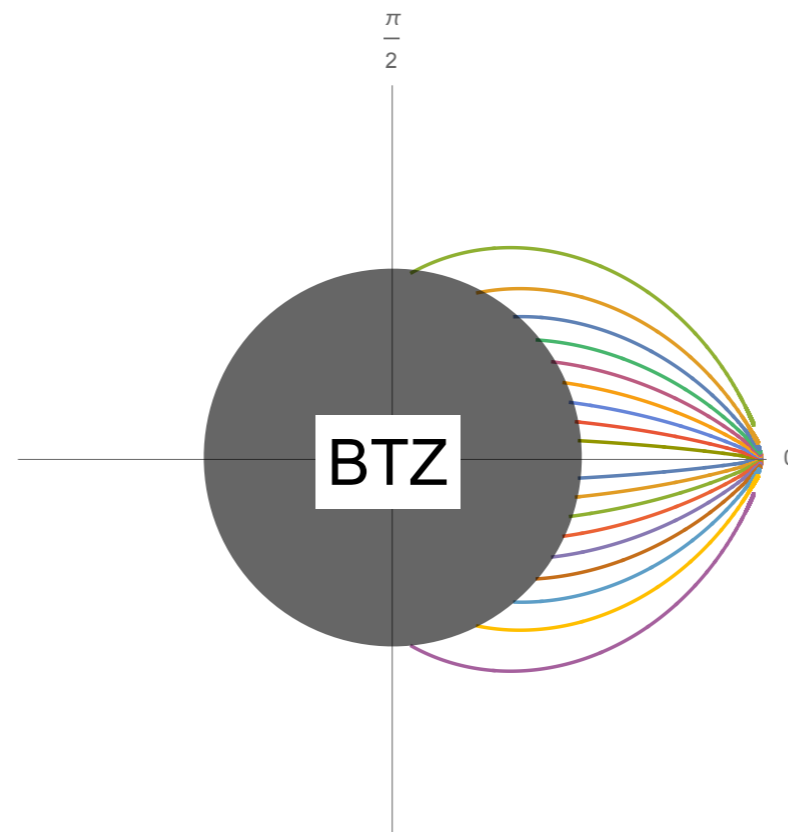


Difficulty: processing material to sphere

Processing materials near QCP is not easy

If only it's ring...

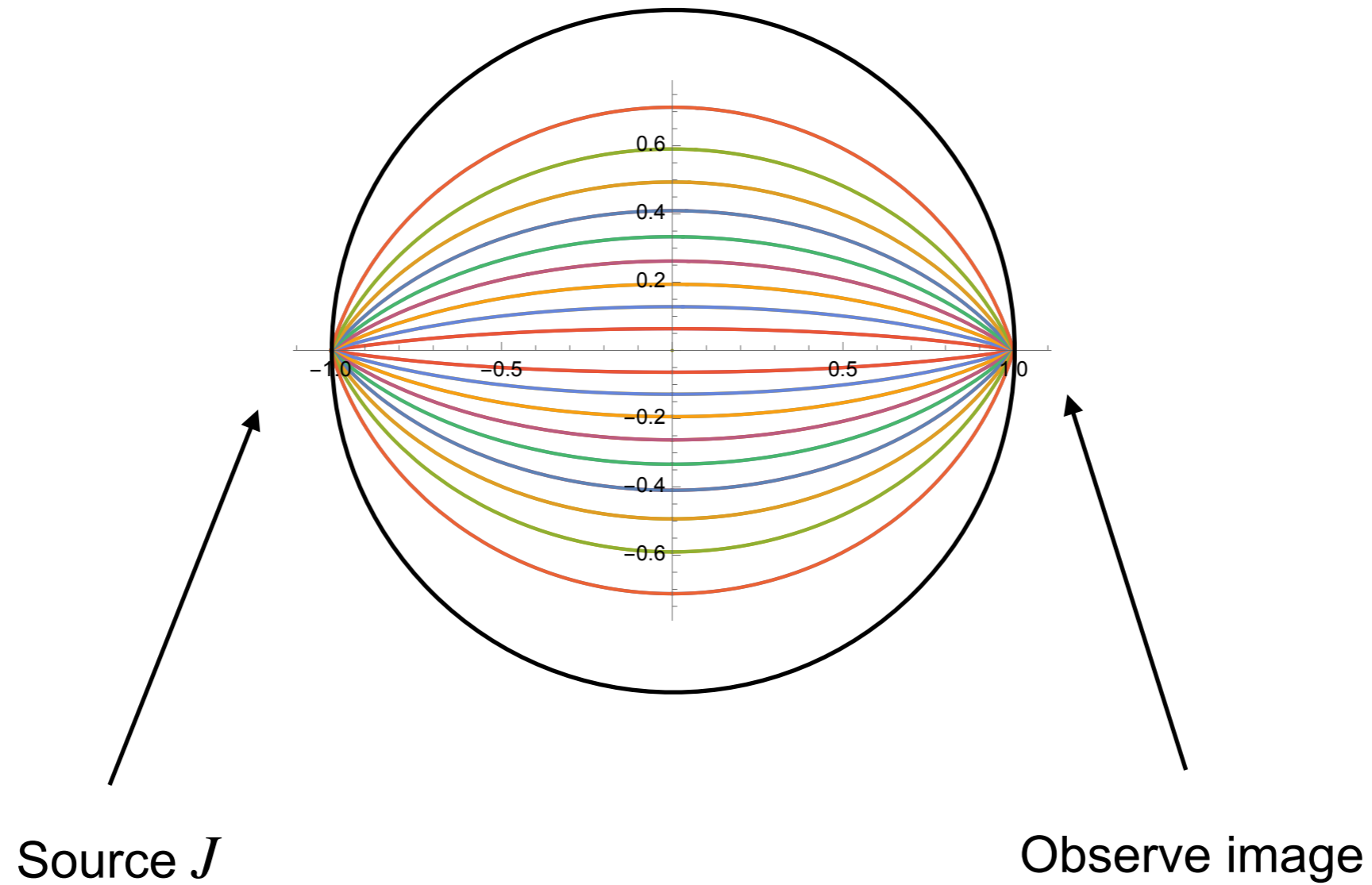
But if it's ring ...



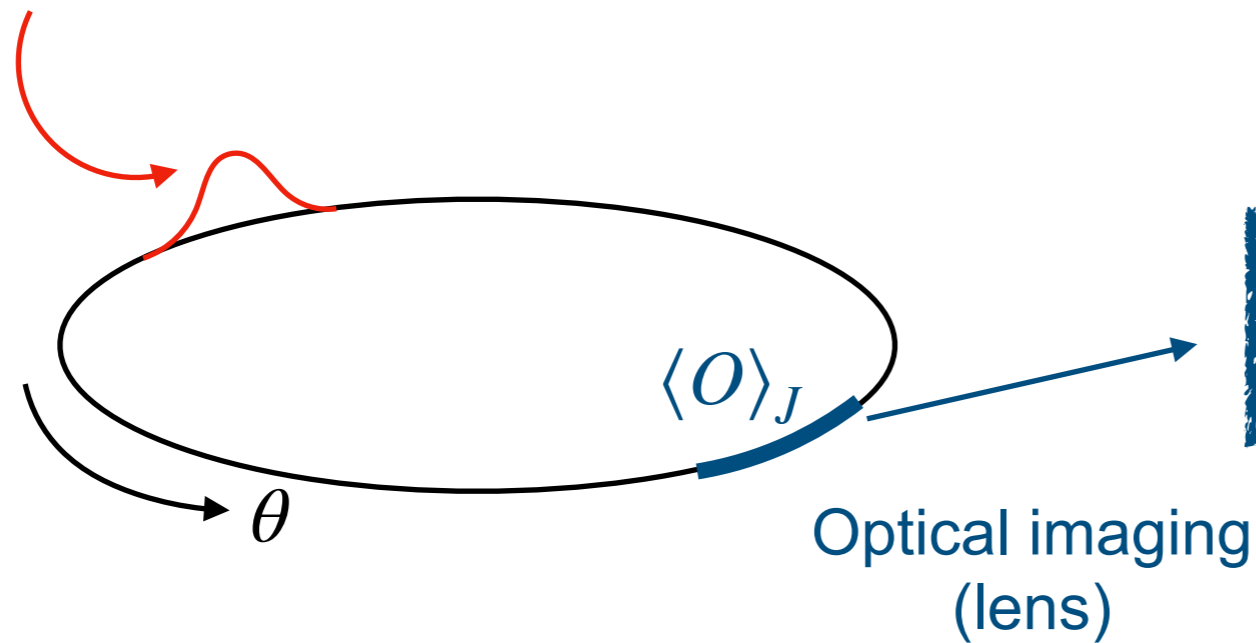
No Einstein ring

Giving up BH, let's find AdS₃

Null geodesics in pure AdS₃

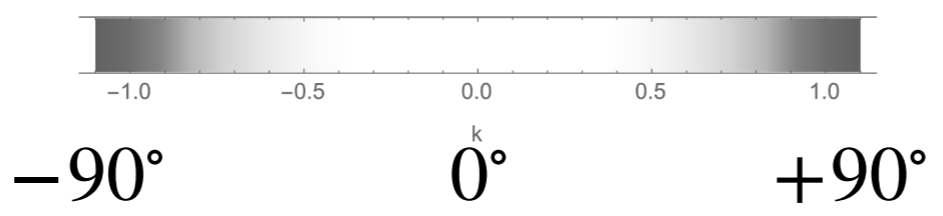


$$J \sim \exp \left[-\frac{\theta^2}{2\sigma^2} - i\omega t \right]$$



$$\psi(t, k) \propto \int_{-d}^d d\theta \langle O(t, \theta) \rangle_J e^{-i\omega k \theta}$$

Your eyesight at $\theta = \pi$



AdS!

Optical imaging reveals spacetime emergence

1. Demonstration with bulk scalar
2. Any bulk field will do
3. Experiment seems technically feasible

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The model: Einstein + probe scalar

$$S = \frac{1}{16\pi G} \int d^3x \sqrt{-g} (R - 2\Lambda) + \int d^3x \sqrt{-g} \left(-\frac{1}{2} \partial_M \Phi \partial^M \Phi \right)$$

Hawking-Page transition

$g_{\mu\nu}$: fixed

$T > \frac{1}{2\pi L}$ $ds^2 = -\frac{r^2 - r_h^2}{L^2} dt^2 + \frac{L^2}{r^2 - r_h^2} dr^2 + r^2 d\theta^2$ $r_h = 2\pi L^2 T$

$T < \frac{1}{2\pi L}$ $ds^2 = -\frac{r^2 + L^2}{L^2} dt^2 + \frac{L^2}{r^2 + L^2} dr^2 + r^2 d\theta^2$

Physics of the material

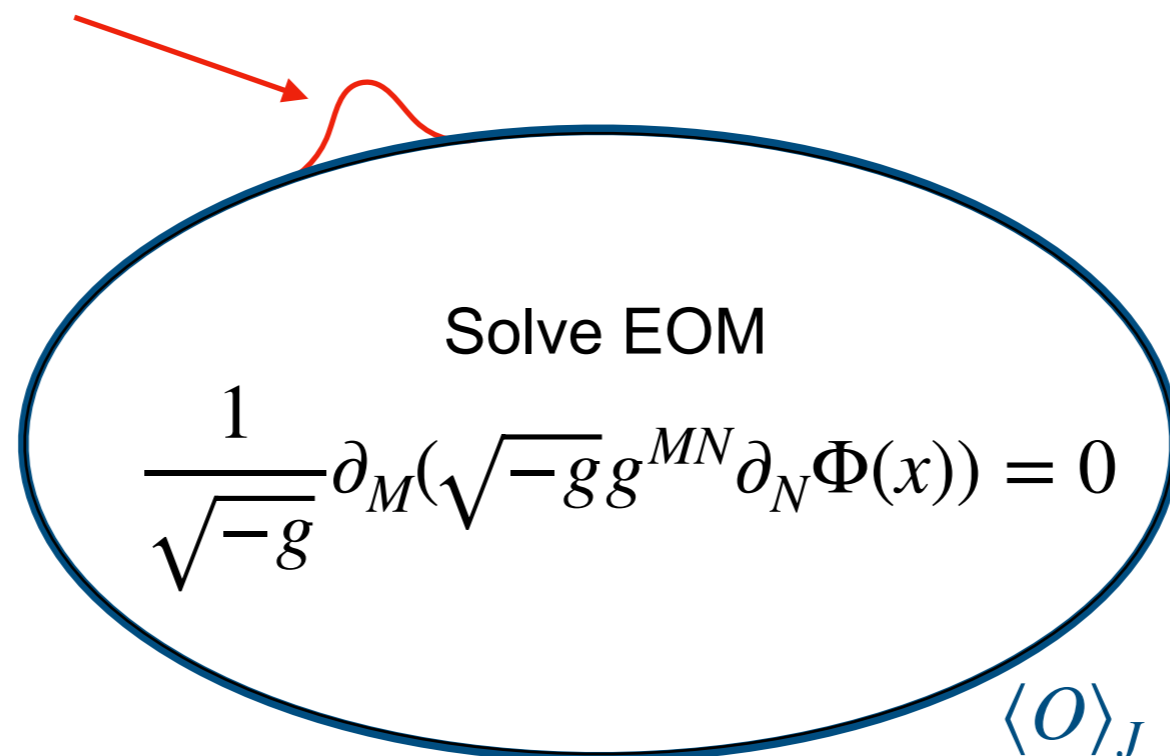
$$\frac{1}{\sqrt{-g}} \partial_M (\sqrt{-g} g^{MN} \partial_N \Phi(x)) = 0 \quad \text{with GKPW dictionary}$$

J is leading, $\langle O \rangle_J$ is subleading

GKPW formula

$$\Phi(x) \sim J(t, \theta) + \frac{\langle O(t, \theta) \rangle_J}{r^2} \quad (r \rightarrow \infty)$$

$$J \sim \exp \left[-\frac{\theta^2}{2\sigma^2} - i\omega t \right] \quad \text{Boundary condition (input)}$$

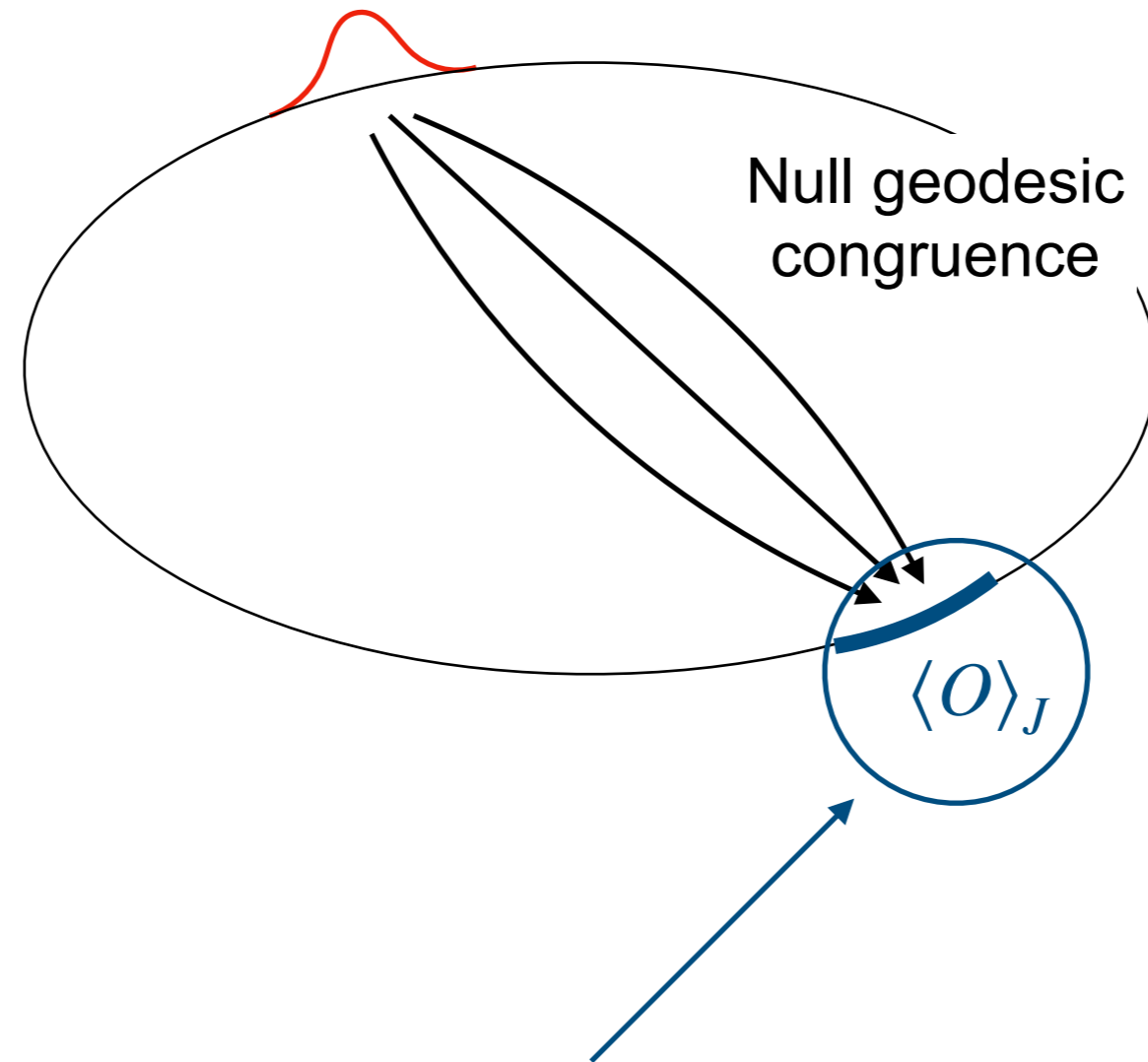


Solve EOM

$$\frac{1}{\sqrt{-g}} \partial_M (\sqrt{-g} g^{MN} \partial_N \Phi(x)) = 0$$

$\langle O \rangle_J$ (output)

Optical imaging recovers image of J

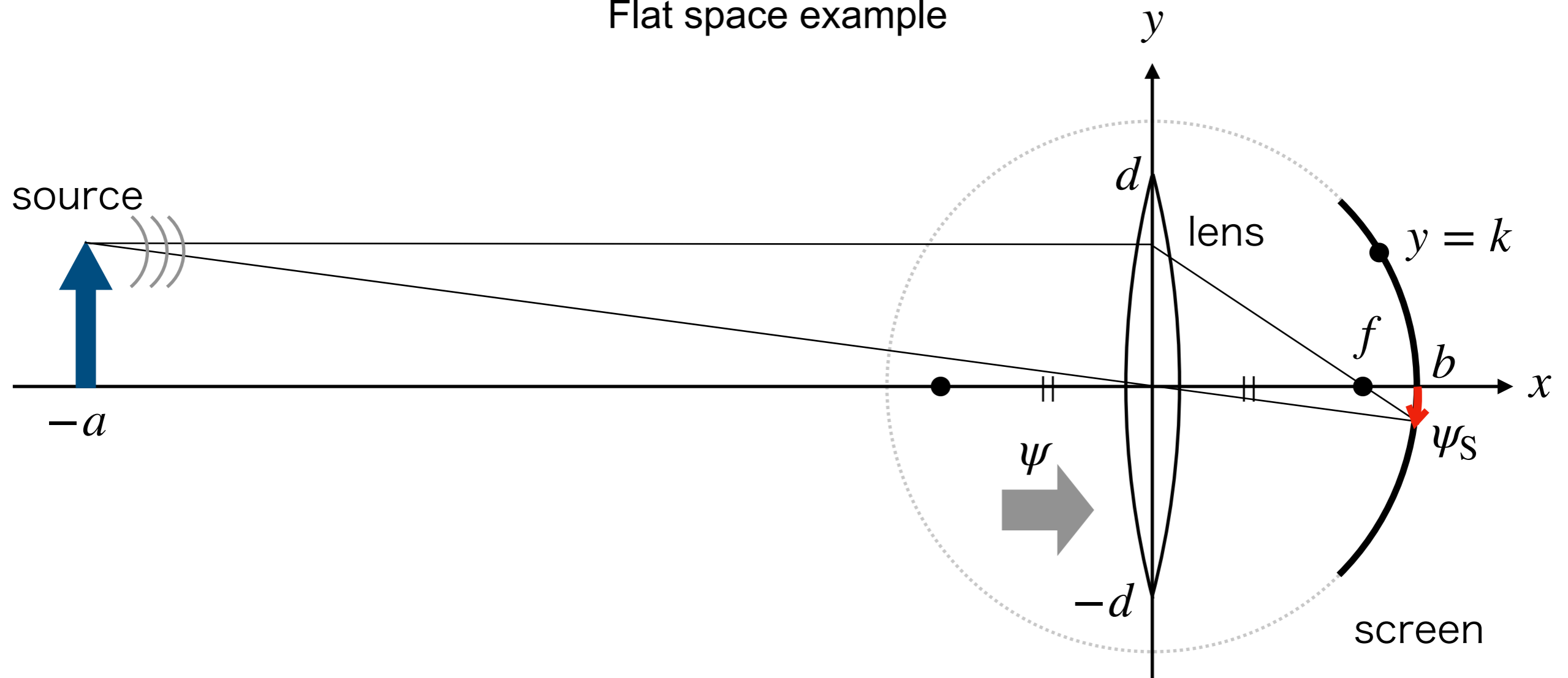


$$\psi(t, k) \propto \int_{-d}^d d\theta \langle O(t, \theta) \rangle_J e^{-i\omega k \theta}$$

Optical imaging
(lens)

Optical imaging recovers image of J

Flat space example

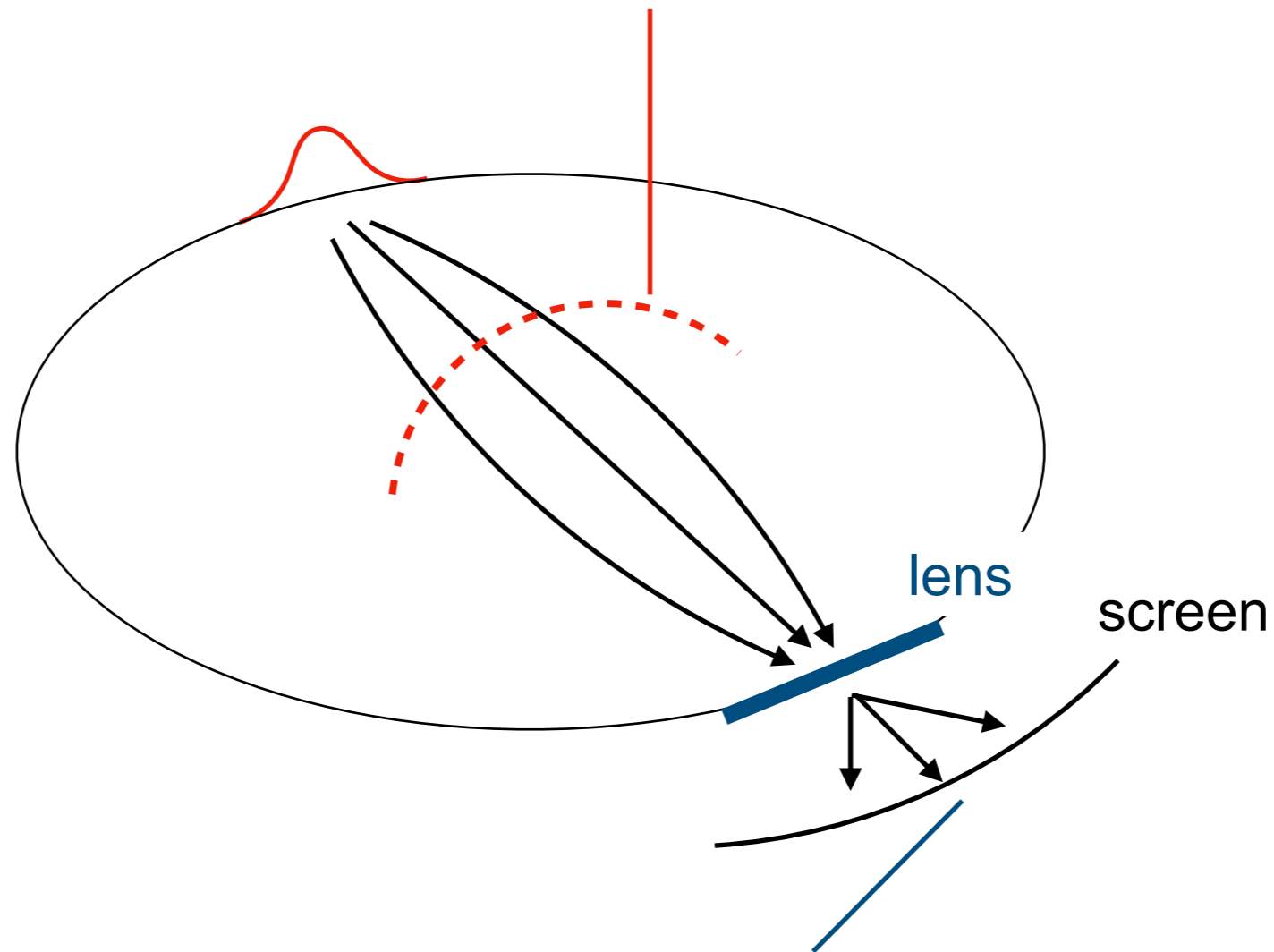


Lens recovers the source image

$$\psi_S(k) \propto \int_{-d}^d dy \psi(y) e^{-i\omega ky/f}$$

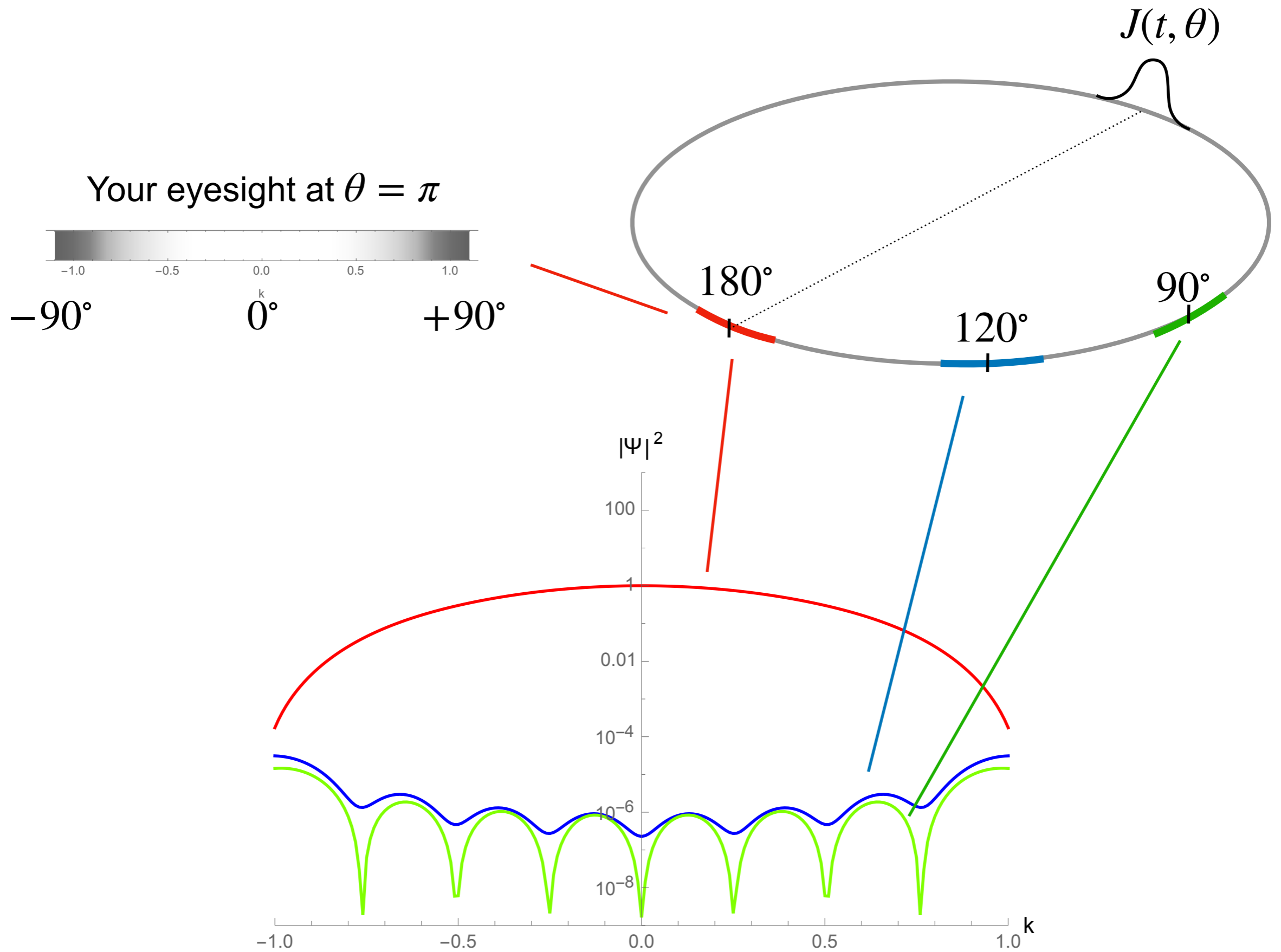
Optical imaging recovers **virtual** image of J

Lens recovers the virtual image of J



$$\psi(t, k) \propto \int_{-d}^d d\theta \langle O(t, \theta) \rangle_J e^{-i\omega k \theta}$$

Plots



Comparison with non-holographic model

Real scalar theory on the ring

$$S = \int d^2x \left(-\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 + J\phi \right)$$

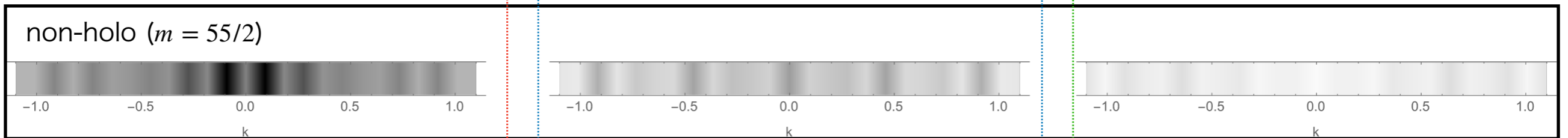
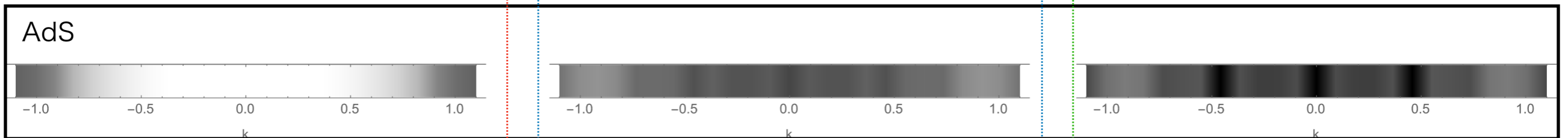
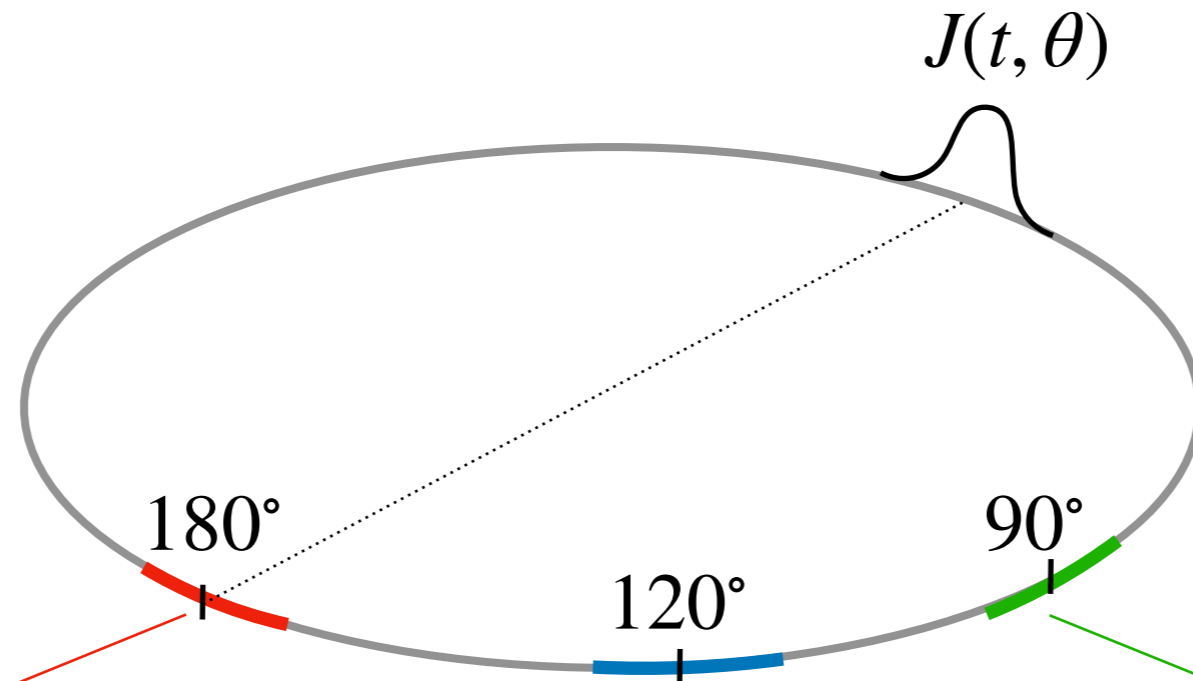
coupling to J

$$J \sim \exp \left[-\frac{\theta^2}{2\sigma^2} - i\omega t \right]$$

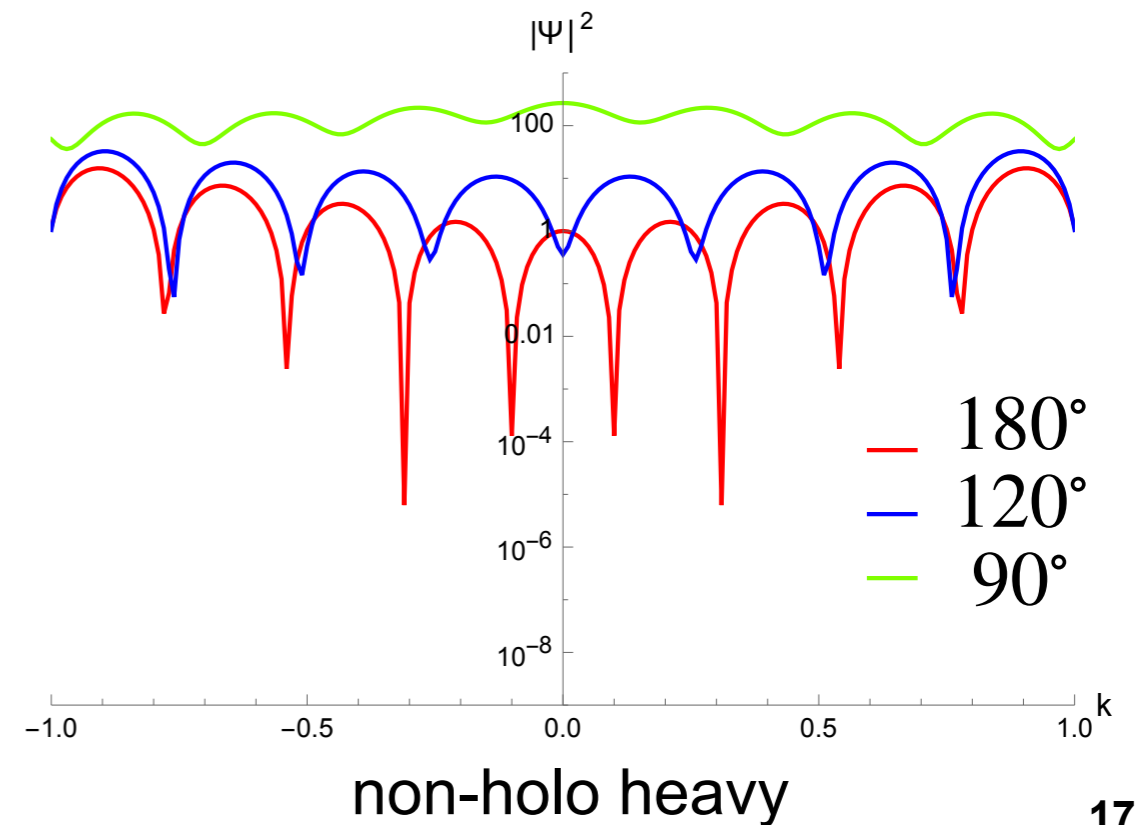
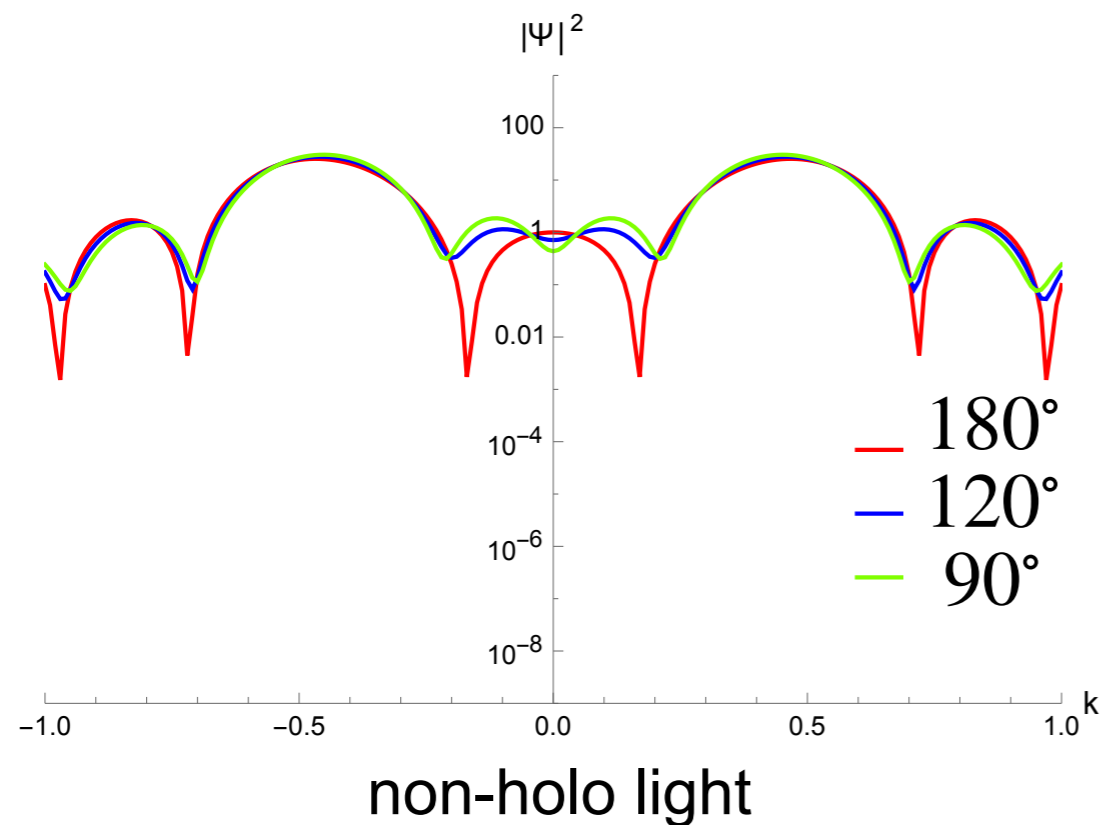
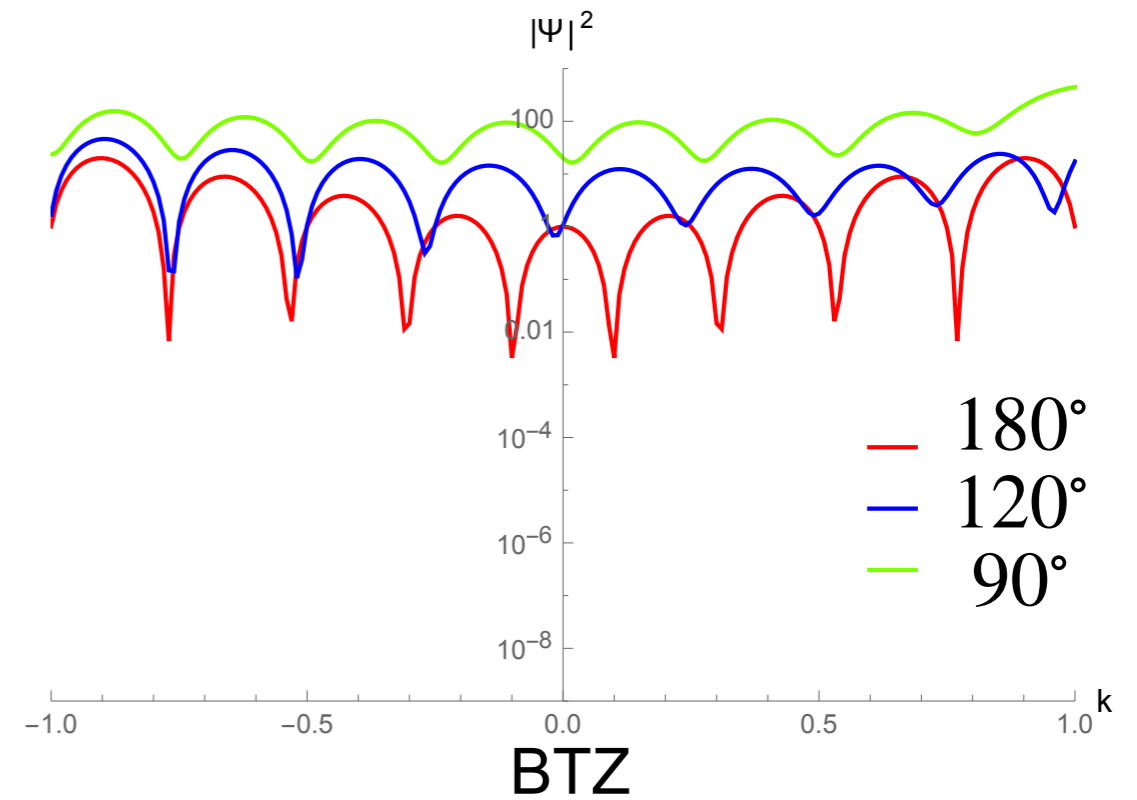
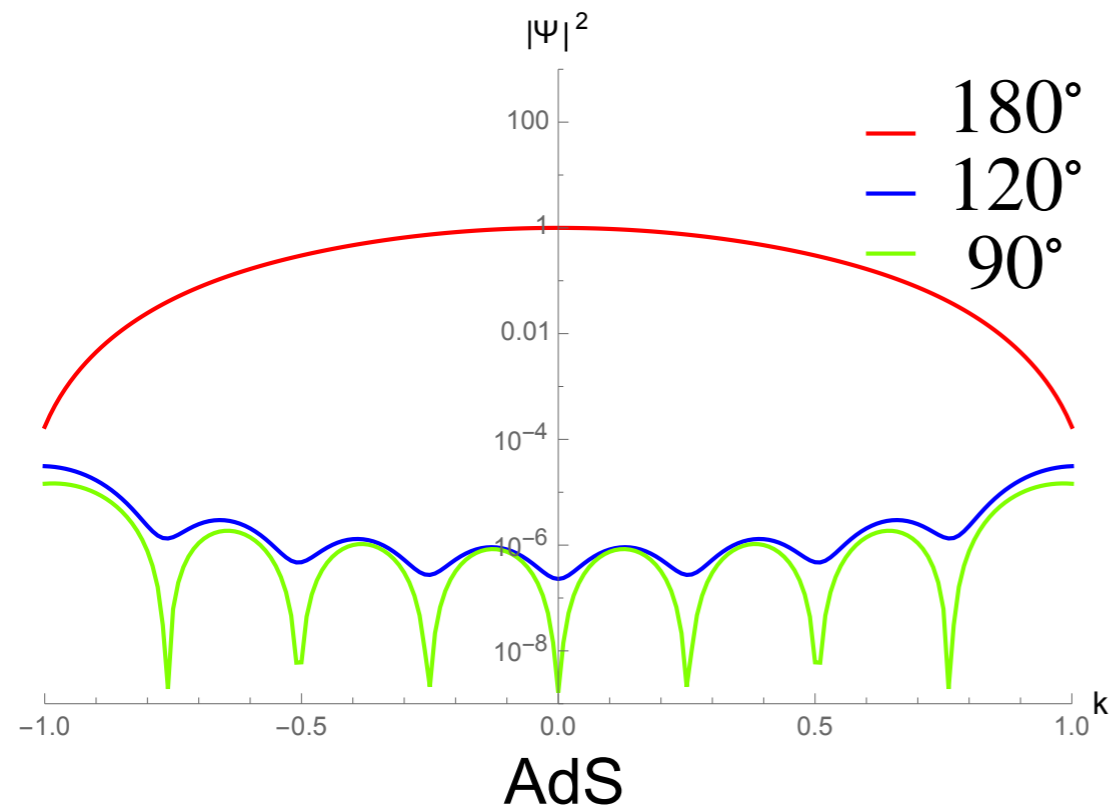
EOM

$$(\partial^2 - m^2)\phi = J$$

Comparison with non-holographic one



Comparison with non-holographic one



Optical imaging reveals spacetime emergence

1. Demonstration with bulk scalar
2. Any bulk field will do
3. Experiment seems technically feasible

Eikonal approx: wave \leftrightarrow congruence

If we add mass?

$$\text{Wave : } \Phi = a(x)e^{iS(x)}$$

$$\begin{aligned} 0 &= (\nabla^\mu \nabla_\mu - m^2)(ae^{iS}) \\ &= \nabla^\mu \left[e^{iS} \left(\nabla_\mu a + ia\partial_\mu S \right) \right] - m^2 ae^{iS} \\ &= e^{iS} \left[\underbrace{-a(\partial_\mu S)^2}_{\text{red circle}} + \cancel{2i\partial^\mu S \nabla_\mu a} + \cancel{\nabla^\mu \nabla_\mu a} + \cancel{ia \nabla^\mu \partial_\mu S} - \cancel{am^2} \right] \end{aligned}$$

Eikonal approx : keep $|\partial_\mu S|$ large

In general, in weak coupling theory,

$$\text{Eikonal approx : } \partial_\mu S \partial^\mu S = 0$$

Eikonal approx: wave \leftrightarrow congruence

$$\partial_\mu S \partial^\mu S = 0$$

$\downarrow \nabla^\nu$

$$\partial_\mu S \nabla^\mu \partial_\nu S = 0$$



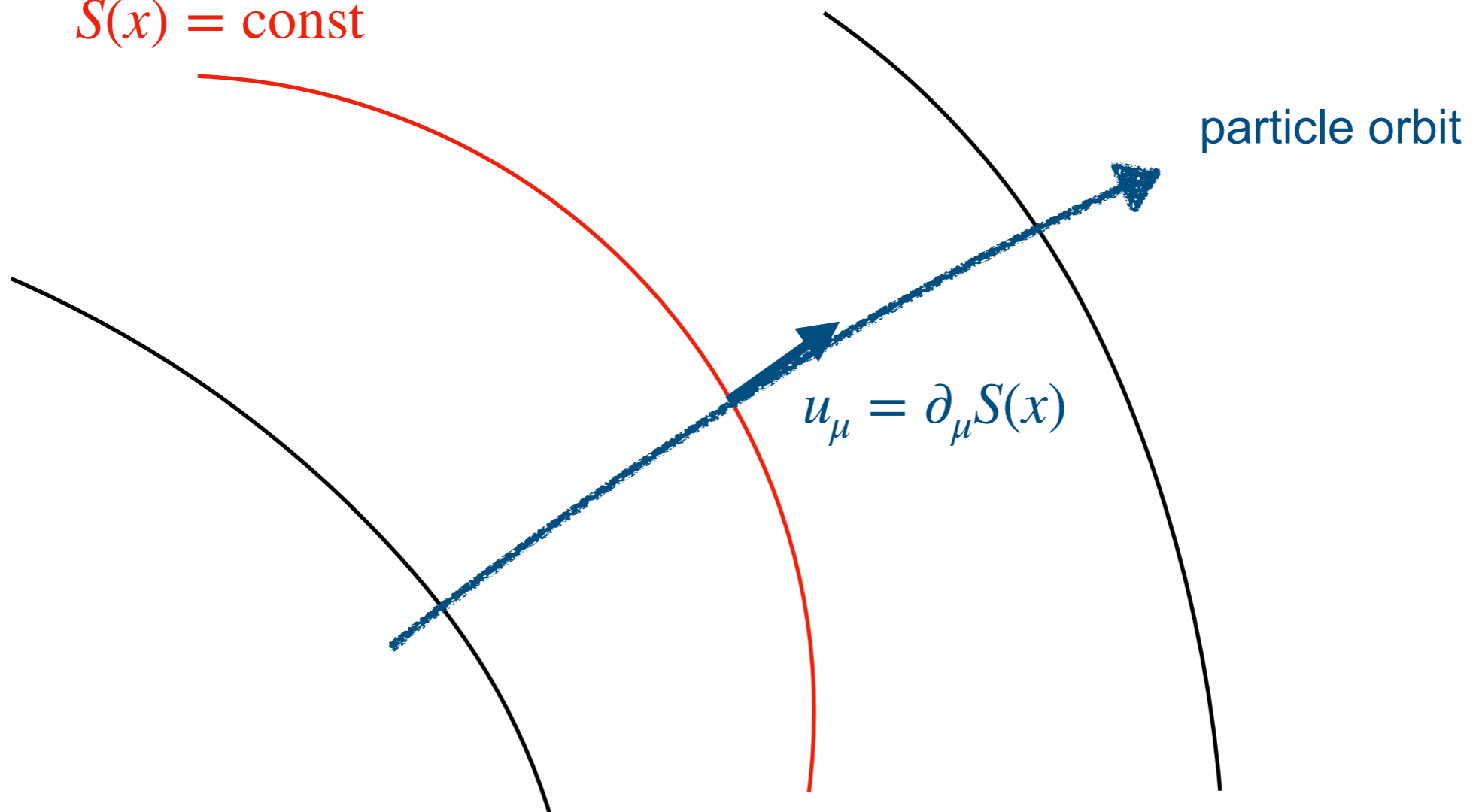
$$u_\mu = \partial_\mu S$$

Geodesic eq

$$u_\mu u^\mu = 0$$

$$\nabla_u u^\mu = 0$$

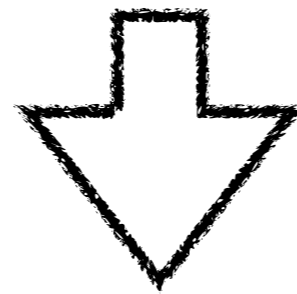
$S(x) = \text{const}$



Any bulk field is available with ω kept large

$$J \sim \exp \left[-\frac{\theta^2}{2\sigma^2} - i\omega t \right]$$

Boundary condition for $S(x)$



As long as ω is kept large, the wave equation is reduced to

$$\partial_\mu S \partial^\mu S = 0$$

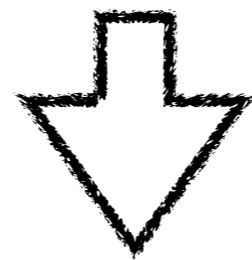
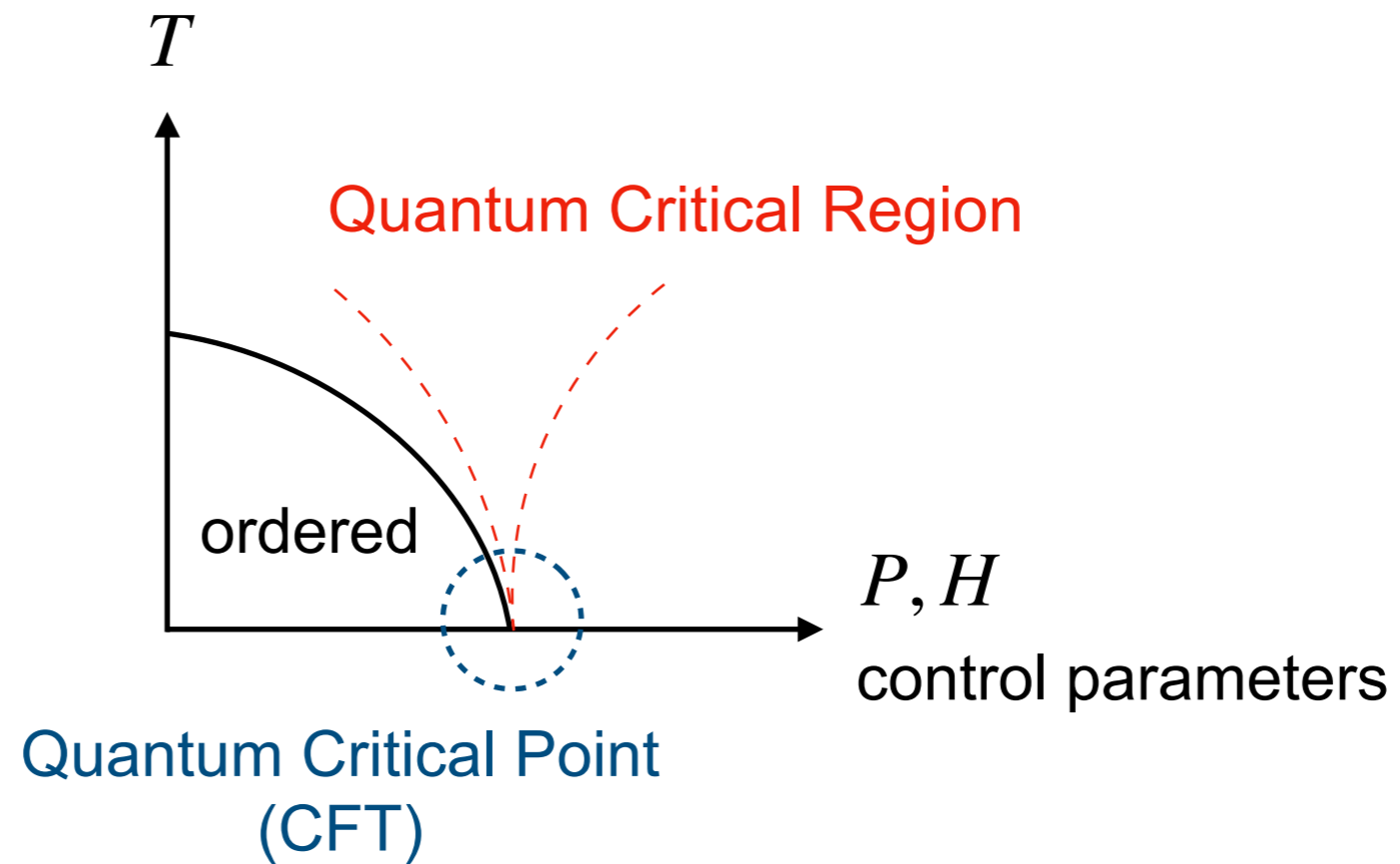
$$\partial_\mu S \nabla^\mu \partial_\nu S = 0$$

Regardless of spin or mass

Optical imaging reveals spacetime emergence

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Probing QCP



Process it to ring

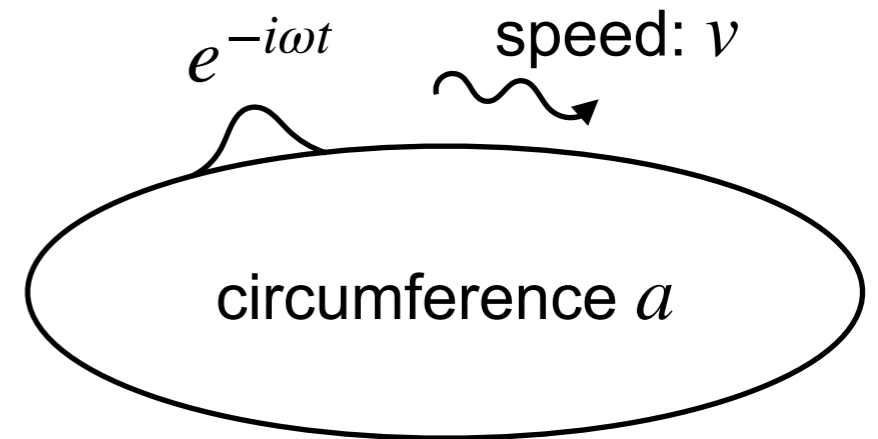
CFT on S^1

Estimation: we can do the experiment

Wavelength \ll ring size $\omega \gg \frac{2\pi v}{a}$ $\omega \sim \frac{50\pi v}{a}$

Low temperature phase $\frac{v\hbar}{a} > k_B T$

Continuum limit is good $a \gg 1 \text{ nm}$

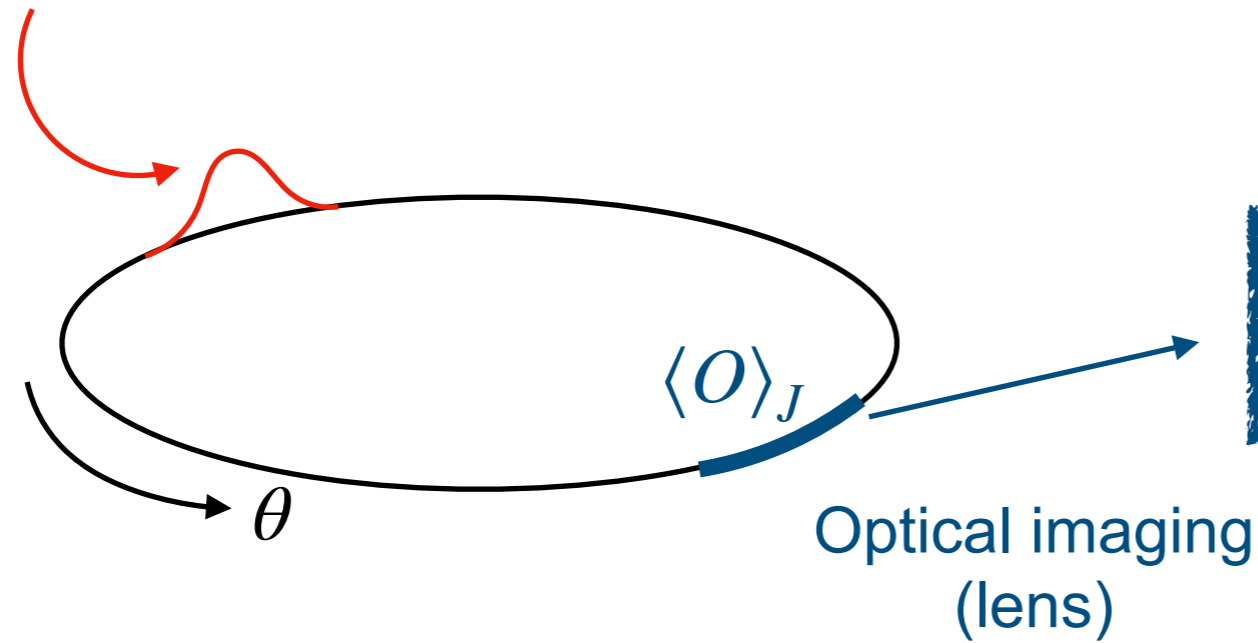


TiCuCl₃ J : EM field, O : spin operator

$$v \sim 2.0 \times 10^3 \text{ m/s}$$

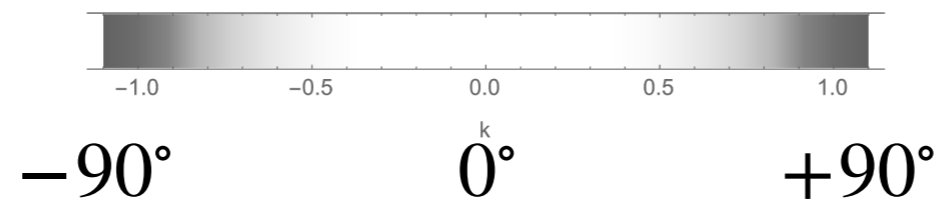
$$T \sim 0.1 \text{ K}, \quad a \sim 10 \text{ nm}$$

$$J \approx \exp \left[-\frac{\theta^2}{2\sigma^2} - i\omega t \right]$$



$$\psi(t, k) \propto \int_{-d}^d d\theta \langle O(t, \theta) \rangle_J e^{-i\omega k \theta}$$

Your eyesight at $\theta = \pi$



AdS!

Response before imaging

