Coarse-graining black holes out of equilibrium with boundary observables on time slice

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Self introduction

Name Daichi Takeda

Status PhD student in Kyoto U (D3 from today)

Field High energy physics

Research topics so far

String field theory

Holography

Bulk metric reconstruction

AdS/CMP

Black hole thermodynamics

Thermodynamics constrains gravity?

Origin of spacetime is unknown

Quantum theory of gravity is necessary

On the other hand, BH is thermodynamic = macroscopic

BH is statistical mechanics of QM?? Strominger-Vafa (1996), recently Kawai-Yokokura

BH thermodynamics will give macroscopic clues to QG!



- 1. BH thermodynamics and problems
- 2. Coarse-grained entropy and 2nd law in CFT
- 3. Rewrite in gravity by AdS/CFT
- 4. 2nd law implies null energy condition
- 5. (Generalized) 1st law in GR

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4 laws in thermodynamics

Oth law: Existence of intensive variables

Temperature T, chemical potential μ , ...

1st law: Energy conservation $dE = TdS + \mu dN + \cdots$

2nd law: Iff $X \to Y$ is possible with X in equilibrium, $S_X \leq S_Y$

3rd law: Entropy vanishes at
$$T = 0$$

Unnecessary to construct thermodynamics

BHT is almost parallel



Oth law: intensive variables T, Ω, \dots

1st law: Energy conservation $dM = TdS + \Omega dP + \cdots$

2nd law: still under debate

Hawking area law? Generalized second law? Other candidates?

2nd law is under debate

Generalized entropy?



I think, this claim is different from thermodynamic 2nd law... 2nd law: If $X \to Y$ is possible, $S_X \le S_Y$, with X in equilibrium



always out of equilibrium

 $S_{\text{gen}} \neq S_{\text{Sch}}$ for initial state

To do: Find entropy obeying 1st and 2nd law



from matter fields

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Coarse-graining is respecting some aspects

Canonical ensemble

Maximize $S(\rho) = -\operatorname{Tr} \rho \ln \rho$ under $\operatorname{Tr}(\rho H) = E$ and $\operatorname{Tr} \rho = 1$



$$S_{\rm can} = -\,{\rm Tr}\,\rho_{\rm can}\,{\rm ln}\,\rho_{\rm can}$$

Coarse-graining is respecting some aspects

Coarse-grained state $ho_{
m cg}$

 $\{H, P_A, O_I(\theta)\}$: operator set to be respected

 $\begin{array}{l} \text{Maximize } S(\rho) = -\operatorname{Tr}\rho\ln\rho\\ \text{under } \mathrm{Tr}(\rho H) = h, \ \mathrm{Tr}(\rho P_A) = p_A, \ \mathrm{Tr}(\rho O_I(\theta)) = o_I(\theta), \ \mathrm{Tr}\rho = 1 \end{array}$

$$\rho_{\rm cg} = \frac{1}{Z} \exp\left[-\beta \left(H - \omega^A P_A - \int d^{d-1}\theta \,\lambda^I(\theta) O_I(\theta)\right)\right]$$

Coarse-grained entropy

$$S := -\operatorname{Tr}\rho_{\rm cg}\ln\rho_{\rm cg}$$

Coarse-grained entropy of time *t*

Coarse-graining conditions: $Tr(\rho H) = h$, $Tr(\rho P_A) = p_A$, $Tr(\rho O_I(\theta)) = o_I(\theta)$



2nd law from relative entropy



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AdS/CFT constrains BH dynamics



Setup: equilibrium to non-equilibrium



GKPW formula and 1pt functions





 $\left\langle e^{i\int \mathrm{d}^{d-1}\theta j^{I}(t,\theta)O_{I}(\theta)}\right\rangle$ $e^{iI_{\text{grav}}[\Phi]}$ with $\Phi(x) \sim r^{\Delta_I - d} j^I(t, \theta)$ $\frac{\delta}{\delta j^{I}(t,\theta)} I_{\text{grav}}[\Phi] =: \pi_{I,t}(\theta)$ $\operatorname{Tr}(\rho_t O_I(\theta))$ $\operatorname{Tr}(\rho_t H), \operatorname{Tr}(\rho_t P_A)$ ADM mass, momenta

=

Computed from Brown-York tensor

Coarse-grained state = Euclid BH



At each time t, coarse-graining conditions become

 $\pi_{I,t}(\theta) = \pi_{I}^{(E)}(\theta)$, and matching of mass and momenta

Solution:
$$(\beta, \omega, \lambda) = (\beta_t, \omega_t, \lambda_t)$$

Entropy is the cigar tip area



 $A_t \ge A_0$ via AdS/CFT





 $A_t \ge A_0$

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Sch-AdS



The values to be respected: mass M_t , angular momenta $P_{A,t}$, charge Q_t

$$M_t = \frac{d-1}{8\pi G} \text{Vol}(\mathbb{S}^{d-1}) \times \mu(t) + (\mu - \text{indep.})$$
$$P_{A,t} = 0$$
$$Q_t = 0$$

Sch-AdS



Euclid BH having same $M_t, P_{A,t}, Q_t$

$$ds^{2} = f(v, r)d\tau^{2} + \frac{dr^{2}}{f(v, r)} + r^{2}d\Omega^{2}$$
$$f(v, r) = 1 + \frac{r^{2}}{L^{2}} - \frac{2\mu(v)}{r^{d-2}}$$



AdS/CFT says $A_t \geq A_0$ This does not hold for all $\mu(v)$



The same thing holds for other cases

 $\label{eq:AdS/CFT} \begin{array}{l} \text{AdS/CFT says} \\ A_t \geq A_0 \\ \\ \text{This does not hold always} \end{array}$

But it holds if $T_{\ell\ell} \geq 0$

I also confirmed in $P_{\phi} \neq 0$ case and $Q \neq 0$ case

asymptotically flat

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Entropy vs Euclid action: Legendre tr

Entropy
$$S_t$$
 and free energy $F_t := \beta_t^{-1} I_{\text{grav}}^{(\text{E})}[\beta_t, \omega_t, \lambda_t]$

$$S_{t} = -I_{\text{grav}}^{(\text{E})}[\beta_{t}, \omega_{t}, \lambda_{t}] + \beta_{t} \left(M_{t} - \omega_{t}^{A} P_{A,t} - \int d^{d-1} \vec{x} \, \lambda_{t}^{I}(\theta) \pi_{I,t}(\theta) \right)$$

$$Values to be respected$$

 $(\beta_t, \omega_t, \lambda_t)$ are viewed as functions of (M_t, P_t, π_t)

c.f.) CFT description

$$S_{t} = -\operatorname{Tr}\rho_{\mathrm{cg},t}\ln\rho_{\mathrm{cg},t}$$
$$= \ln Z[\beta_{t},\omega_{t},\lambda_{t}] + \beta_{t}\left(\langle H\rangle_{t} - \omega^{A}\langle P_{A}\rangle_{t} - \int \mathrm{d}^{d-1}\theta\,\lambda^{I}(\theta)\langle O_{I}(\theta)\rangle_{t}\right)$$

First law is generalized

$$S_t = -I_{\text{grav}}^{(\text{E})}[\beta_t, \omega_t, \lambda_t] + \beta_t \left(M_t - \omega_t^A P_{A,t} - \int d^{d-1}\theta \,\lambda_t^I(\theta) \pi_{I,t}(\theta) \right)$$

 $(\beta_t, \omega_t, \lambda_t)$ are viewed as functions of (M_t, P_t, π_t)

The time dependence of S_t is through $(M_t, P_{A,t}, \pi_{I,t})$

The variation of
$$I_{\text{grav}}^{(\text{E})}[\beta, \omega, \lambda]$$

 $\delta I_{\text{grav}}^{(\text{E})}[\beta, \omega, \lambda] = M\delta\beta - P_A\delta(\beta\omega) - \beta \int d^{d-1}\theta \,\delta\lambda^I(\theta)\pi_I(\theta) + (\text{EOM})$

Setting
$$\delta = \frac{\mathrm{d}}{\mathrm{d}t}$$
, we obtain

$$\dot{S}_t = \beta_t (\dot{M}_t - \omega_t^A \dot{P}_{A,t}) - \int d^{d-1}\theta \,\lambda_t^I(\theta) \,\dot{\tilde{\pi}}_{I,t}(\theta), \quad \tilde{\pi}_{I,t} = \beta_t \pi_{T,t}$$

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